



## Mild solutions for conformable fractional order functional evolution equations via Meir-Keeler type fixed point theorem

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**Abstract.** In this study, we delve into the realm of mild solutions for conformable fractional order functional evolution equations, focusing on cases where the fractional order is strictly greater than 1 and less than 2 within a separable Banach space. We demonstrate the existence, uniqueness, attractivity, and controllability of these solutions under local conditions. Our approach involves leveraging a contribution of Meir-Keeler's fixed point theorem alongside the principle of measures of noncompactness. To demonstrate the practical ramifications of our theoretical finds, we provide a specific example that underscores the relevance and applications of the established results.

### 1. Introduction

Fractional calculus extends classical differentiation and integration to non-integer orders, thereby unifying discrete and continuous domains. It commonly employs definitions like the Riemann-Liouville and Caputo formulations (see eg. [40, 44]). Khalil *et al.* [39] proposed the conformable derivative, which shares important properties with the integer-order derivative. The conformable derivative has been applied to many sciences such as physics (see eg. [16, 41, 46]), biology (see eg. [15, 49]), chemistry (see eg. [45]), and medicine (see eg. [11, 13]).

Jaiswal *et al.* [34] studied a conformable fractional abstract initial value problem in Banach spaces, finding moderate solutions using the contraction principle. Kataria *et al.* [37] embarked on an exploration of mild solutions for impulsive integro-differential equations, harnessing the power of conformable differential operators and fixed-point theorems. Among all we may count some of the contributions to conformable fractional evolution equations have also been made by Bouaouid *et al.* [26, 28] and Boukenkoul *et al.* [29].

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This work leverages measures of noncompactness (see eg. [12, 21, 35]) to address solution existence issues and illuminate nonlinear differential equations (see eg. [1, 9, 10, 22, 23, 24, 25, 47]).

Additionally, Baghli *et al.* [20] utilized the nonlinear Leray-Schauder alternative [32] for controllability within bounded intervals, while Agarwal *et al.* [6] addressed controllability over semi-infinite intervals using the nonlinear alternative of Avramescu [19] for equations with infinite delay.

In their study [50], Zou *et al.* established four uniqueness criteria for the initial value problem of differential equations incorporating the conformable fractional derivative. Their work extends Nagumo-type uniqueness theory and Lipschitz conditional theory. To bolster these theoretical advancements, they provide four practical examples, showcasing the efficacy of their approach in proving fractional differential equations.

Bouaouid *et al.* [27] applied Krasnoselskii's fixed point theorem to obtain mild solutions for an abstract class of conformable fractional differential equations. Building on this foundation, Atraoui *et al.* [17] ventured further by utilizing the Darbo-Sadovskii fixed point theorem.

Moreover, Atraoui *et al.* [18] ventured into the intricate realm of controllability, exploring mild solutions for a nonlocal fractional conformable Cauchy problem with keen insight by using second-order differential equations. They used the Banach contraction principle and the cosine linear operator family to arrive at their conclusions.

This work examines conformable fractional order evolutionary equations with local conditions:

$$D^c[D^c\psi(s)] = \mathfrak{F}\psi(s) + \Psi(s, \psi_s), \quad \text{a.e. } s \in I := [0, +\infty); \quad (1)$$

$$\psi(s) = \eta(s), \quad s \in \mathcal{H} := [-a, 0], \quad \text{where } 0 < a < +\infty, \quad D^c\psi(0) = \vartheta \in \mathcal{F}; \quad (2)$$

such that,  $\Psi : I \times C([-a, 0], \mathcal{F}) \rightarrow \mathcal{F}$  denotes a specified function,  $\eta : [-a, 0] \rightarrow \mathcal{F}$  is continuous, and  $\mathfrak{F} : D(\mathfrak{F}) \subset \mathcal{F} \rightarrow \mathcal{F}$  serves as the infinitesimal generator of a strongly continuous cosine function composed of bounded linear operators  $\{\mathfrak{C}(s)\}_{s \in \mathbb{R}}$  such that  $\mathfrak{S}(s) = \int_0^s \mathfrak{C}(x) dx$ , and  $\mathcal{F}$  denotes real separable Banach space generates by the norm  $|\cdot|$ . Furthermore,  $D^c$  is a fractional conformable derivative, where  $0 < c \leq 1$ .

We define  $\psi_s$  for  $s \geq 0$  as a continuous function from  $\mathcal{H}$  to  $\mathcal{F}$  given by:  $\psi_s(x) = \psi(s+x)$ , where  $\psi_s(\cdot)$  represents the state's history from  $s-a$  to the present moment  $s$ .

Subsection 3.3 investigates the attractivity of mild solutions under these settings. Subsection 3.4 delves into how mild solutions can be controlled across the interval  $I = [0, +\infty)$ , providing a thorough exploration of their behavior, so we consider the following problem:

$$D^c[D^c\psi(s)] = \mathfrak{F}\psi(s) + \Psi(s, \psi_s) + \mathcal{B}\mathcal{U}(s), \quad \text{a.e. } s \in I := [0, +\infty); \quad (3)$$

$$\psi(s) = \eta(s), \quad s \in \mathcal{H} := [-a, 0], \quad \text{where } 0 < a < +\infty, \quad D^c\psi(0) = \vartheta \in \mathcal{F}; \quad (4)$$

where  $\mathfrak{F}$ ,  $\Psi$ ,  $\eta$  and  $\vartheta$  are as in problem (1) – (2), the control function  $\mathcal{U}(\cdot)$  finds its place within  $L^2(I, \mathcal{F})$ , the Banach space housing admissible control functions, while  $\mathcal{B}$  represents a bounded linear operator mapping from  $\mathcal{F}$  to  $\mathcal{F}$ .

Ultimately, we provide an illustrative example demonstrating the abstract theory expounded in the preceding sections.

## 2. Preliminary Concepts

In this section, we present symbols, explanations, and fundamental principles drawn from multivalued analysis. These elements will be incorporated consistently in the subsequent sections of this paper.

Let  $C(\mathcal{H}, \mathcal{F})$  be the Banach space of continuous functions with the norm

$$\|x\| = \sup\{|x(s)| : s \in \mathcal{H}\}.$$

$BC(I, \mathcal{F})$  denotes the Banach space comprising all functions from  $I$  to  $\mathcal{F}$  that are both bounded and

continuous, endowed with the norm

$$\|x\|_{BC} = \sup\{|x(s)| : s \in I\}.$$

Consider the space  $BC_\infty$  defined as  $\{x : [-a, +\infty) \rightarrow \mathcal{F}, x|_{[0,b]}$  is bounded and continuous for  $b > 0\}$ , with the norm

$$\|x\|_{BC_\infty} = \sup\{|x(s)| : s \in [0, T]\},$$

where  $T = \sup\{b > 0 : x|_{[0,b]}$  is bounded and continuous $\}$ .

**Definition 2.1.** (Khalil et al. [39])

The conformable fractional derivative of order  $0 < c \leq 1$  for a function  $x(\cdot)$  is expressed as

$$D^c x(s) = \lim_{\tau \rightarrow 0} \frac{x(s + \tau s^{1-c}) - x(s)}{\tau}, \quad s > 0;$$

$$D^c x(0) = \lim_{\tau \rightarrow 0} D^c x(\tau).$$

Additionally, the conformable fractional integral of order  $c$  of a function  $x$  is given by

$$I^c x(s) = \int_0^s t^{c-1} x(t) dt$$

provided that the limits and integrals are well defined.

**Definition 2.2.** A map  $f : J \times E \rightarrow E$  is said to be Carathéodory if it satisfies :

- (i)  $y \mapsto f(t, y)$  is continuous for almost all  $t \in J$ ;
- (ii)  $t \mapsto f(t, y)$  is measurable for each  $y \in E$ .

Let us now review some key aspects of the Kuratowski noncompactness measure.

**Definition 2.3.** (see eg. [12, 21, 35]) Let  $\mathcal{D}_{\mathcal{F}}$  be the bounded subsets of  $\mathcal{F}$ . The map  $\mathfrak{K} : \mathcal{D}_{\mathcal{F}} \rightarrow [0, +\infty)$  is called Kuratowski's non-compactness measure and is defined as follows:

$$\mathfrak{K}(\mathcal{E}) = \inf\{\alpha > 0 : \mathcal{E} \subseteq \bigcup_{j=1}^k \mathcal{E}_j \text{ and } \text{diam}(\mathcal{E}_j) \leq \alpha\}, \text{ here } \mathcal{E} \in \mathcal{D}_{\mathcal{F}}.$$

For properties and more details about the Kuratowski measure of noncompactness (see eg. [12, 21, 35]).

In this paper, we shall solve our problem by using the method of fixed point theory (see e.g. [7, 36, 8]) It belongs to us, inspired by [8], we define the following improved form of the Meir-Keeler condensing operator.

**Definition 2.4.** Let  $\mathcal{E}$  be a nonempty subset of the Banach space  $\mathcal{F}$ . Define  $\mathfrak{K}$  as an arbitrary measure of noncompactness on  $\mathcal{F}$ . We introduce  $\mathfrak{N} : \mathcal{E} \rightarrow \mathcal{F}$  as Meir-Keeler condensing operator if it meets the following criteria:  $\mathfrak{N}$  is both continuous and bounded, and for every  $\beta > 0$ , there exists  $\mu > 0$  such that whenever  $\beta < \mathfrak{K}(\mathcal{R}) < \beta + \mu$ , it follows that  $\mathfrak{K}(\mathfrak{N}(\mathcal{R})) \leq \beta$  is true for any subset  $\mathcal{R}$  which is bounded of  $\mathcal{E}$ .

**Lemma 2.5.** (see [33]) Consider  $\mathcal{F}$  as a Banach space, and let  $\mathcal{E} \subset C(I, \mathcal{F})$  that is both bounded and equicontinuous. Consequently, the map  $\mathfrak{K}(\mathcal{E}(s))$  remains continuous over the interval  $I$ , and  $\mathfrak{K}_I(\mathcal{E})$  equals the supremum value of  $\mathfrak{K}(\mathcal{E}(s))$  for  $s$  in  $I$ .

**Theorem 2.6.** (Meir-Keeler's theorem see [8]) Let  $\mathcal{E}$  be a nonempty, bounded, closed, and convex subset of a Banach space  $\mathcal{F}$ . If  $\mathfrak{N} : \mathcal{E} \rightarrow \mathcal{E}$  is a continuous Meir-Keeler condensing operator, then  $\mathfrak{N}$  guarantees at least one fixed point, and the collection of all such fixed points within  $\mathcal{E}$  forms a compact set.

**Definition 2.7.** (see [43]) A function  $T : \mathbb{R}_+ \rightarrow \mathbb{R}_+$  is called an  $L$ -function if  $T(0) = 0$ ,  $T(s) > 0$  for  $s \in \mathbb{R}_+^*$  and for every  $s \in \mathbb{R}_+^*$  there exists  $\delta > 0$  such that  $T(t) \leq s$  for all  $t \in [s, s + \delta]$ .

**Corollary 2.8.** Let  $\mathcal{E}$  be a nonempty, bounded, closed, and convex subset of a Banach space  $\mathcal{F}$ . If  $\mathfrak{N} : \mathcal{E} \rightarrow \mathcal{E}$  is a continuous operator such that

$$\mathfrak{H}(\mathfrak{N}(\mathcal{R})) \leq T(\mathfrak{H}(\mathcal{R})).$$

For every subset  $\mathcal{R}$  of  $\mathcal{E}$ , where  $\mathfrak{H}$  is an arbitrary measure of noncompactness and  $T$  is an  $L$ -function, it follows that  $\mathfrak{N}$  has at least one fixed point. Furthermore, the collection of all fixed points of  $\mathfrak{N}$  in  $\mathcal{E}$  is compact.

**Proof.** To prove Corollary 2.7, it is sufficient to demonstrate that the operator  $\mathfrak{N}$  is a Meir-Keeler condensing operator, following steps similar to those used in Theorem 2.10 [8]. We will then conclude the proof by applying Theorem 2.5.

**Definition 2.9.** (see [30]) We characterize solutions of equations (1) – (2) as locally attractive if there exists a closed ball  $\overline{B}(\psi^*, \sigma)$  in the space  $BC$ , centered at  $\psi^*$  and with radius  $\sigma$ , such that for any solutions  $\psi$  and  $\tilde{\psi}$  of Equations (1) – (2) within this ball, the following convergence condition holds:

$$\lim_{s \rightarrow +\infty} (\psi(s) - \tilde{\psi}(s)) = 0.$$

### 3. Main Results

#### 3.1. Existence results

In this section, we outline our main result concerning the existence of solutions for problem (1) – (2). Before delving into the details and proof of this result, we first introduce the concept of a mild solution.

**Definition 3.1.** We define the mild solution  $\psi \in C([-a, +\infty), \mathcal{F})$  of the problem (1) – (2) as follows

$$\psi(s) = \begin{cases} \eta(s), & \text{if } s \in \mathcal{H}; \\ \mathfrak{C}\left(\frac{s^c}{c}\right) \eta(0) + \mathfrak{S}\left(\frac{s^c}{c}\right) \vartheta + \int_0^s x^{c-1} \mathfrak{S}\left(\frac{s^c-x^c}{c}\right) \Psi(x, \psi_x) dx, & \text{if } s \in I; \end{cases}$$

We must introduce the following hypotheses, which will be utilized later:

- (i) The function  $\Psi : I \times C(\mathcal{H}, \mathcal{F}) \rightarrow \mathcal{F}$  is carathéodory and there exists a continuous function  $\mathcal{O} : I \rightarrow I$  that satisfies:

$$|\Psi(s, u)| \leq \mathcal{O}(s) \|u\|,$$

$$\mathfrak{H}(\Psi(s, \mathcal{D})) \leq \mathcal{O}(s) \mathfrak{H}(\mathcal{D}),$$

and  $\mathcal{O}^* := \sup_{s \in I} \int_0^s x^{c-1} \mathcal{O}(x) dx < \infty$ , for all  $s \in I$ ,  $u \in C(\mathcal{H}, \mathcal{F})$ , bounded set  $\mathcal{D} \subset C(\mathcal{H}, \mathcal{F})$  and  $0 < c \leq 1$ ;

- (ii) The cosine operator  $\mathfrak{C}(s)_{s \in \mathbb{R}}$  is uniformly continuous and there exist constants  $\mathcal{M}_c^\mathfrak{C}$ ,  $\mathcal{M}_c^\mathfrak{S}$  both greater than zero, such that

$$\sup_{s \in I} \left\| \mathfrak{C}\left(\frac{s^c}{c}\right) \right\| \leq \mathcal{M}_c^\mathfrak{C} \text{ and } \sup_{s \in I} \left\| \mathfrak{S}\left(\frac{s^c}{c}\right) \right\| \leq \mathcal{M}_c^\mathfrak{S}.$$

**Theorem 3.2.** Under the assumptions (i) – (ii) and if  $\mathcal{M}_c^\mathfrak{S} \mathcal{O}^* < 1$ , then the system (1) – (2) is guaranteed to have a mild solution within the space  $BC([-a, +\infty), \mathcal{F})$ .

**Proof.** We reformulate the problem (1) – (2) as fixed-point issue. Let us define the operator  $\mathfrak{N} : BC([-a, +\infty), \mathcal{F}) \rightarrow BC([-a, +\infty), \mathcal{F})$  as follows:

$$\mathfrak{N}(\psi)(s) = \begin{cases} \eta(s), & \text{if } s \in \mathcal{H}; \\ \mathfrak{C}\left(\frac{s^c}{c}\right) \eta(0) + \mathfrak{S}\left(\frac{s^c}{c}\right) \vartheta + \int_0^s x^{c-1} \mathfrak{S}\left(\frac{s^c-x^c}{c}\right) \Psi(x, \psi_x) dx, & \text{if } s \in [0, +\infty). \end{cases}$$

The operator  $\mathfrak{N}$  maps  $BC([-a, +\infty), \mathcal{F})$  into  $BC([-a, +\infty), \mathcal{F})$ . Specifically, for  $\psi \in BC([-a, +\infty), \mathcal{F})$  and for any  $s \in I$  we have:

$$\begin{aligned} |\mathfrak{N}(\psi)(s)| &\leq \|\mathfrak{C}\left(\frac{s^c}{c}\right)\| \|\eta(0)\| + \|\mathfrak{S}\left(\frac{s^c}{c}\right)\| \|\vartheta\| + \int_0^s x^{c-1} \|\mathfrak{S}\left(\frac{s^c-x^c}{c}\right)\| \|\Psi(x, \psi_x)\| dx \\ &\leq \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \int_0^s x^{c-1} \mathcal{O}(x) \|\psi_x\| dx \\ &\leq \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* \|\psi\|_{BC}. \end{aligned}$$

So,  $\mathfrak{N} \in BC([-a, +\infty), \mathcal{F})$ . Furthermore, suppose  $l \geq \frac{\mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\|}{1 - \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^*}$ , and let  $B_l$  denote the closed ball in  $BC([-a, +\infty), \mathcal{F})$  centered at the origin with radius  $l$ . consider  $\psi \in B_l$  and  $s \in I$ , we get

$$|\mathfrak{N}(\psi)(s)| \leq \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* l.$$

Thus, we get that

$$\|\mathfrak{N}(\psi)\|_{BC} \leq l.$$

We embark on the validation process to ensure that  $\mathfrak{N} : B_l \rightarrow B_l$  satisfies the criteria outlined in Meir Keeler’s fixed-point Theorem 2.6.

Firstly, we ascertain the continuity of  $\mathfrak{N}$  within  $B_l$ . Let  $\{\psi_n\}$  be a sequence such that  $\psi_n \rightarrow \psi$  in  $B_l$ . We observe that:

$$|\mathfrak{N}(\psi_n)(s) - \mathfrak{N}(\psi)(s)| \leq \mathcal{M}_c^{\mathfrak{S}} \int_0^s x^{c-1} |\Psi(x, (\psi_x)_n) - \Psi(x, \psi_x)| dx.$$

According to (i), we have  $\Psi(x, (\psi_x)_n) \rightarrow \Psi(x, \psi_x)$  as  $n \rightarrow +\infty$  for almost every  $s \in I$ . Application of the Theorem of convergence dominated by Lebesgue yields:

$$\|\mathfrak{N}(\psi_n) - \mathfrak{N}(\psi)\|_{BC} \rightarrow 0, \text{ as } n \rightarrow \infty.$$

Consequently,  $\mathfrak{N}$  exhibits continuity.

So, by the previous steps we can thus infer that  $\mathfrak{N}(B_l) \subset B_l$ .

Next, we confirm the equicontinuity of  $\mathfrak{N}(B_l)$  over every compact interval  $J$  of  $[0, +\infty)$ . Let  $x_1, x_2 \in J$  with  $x_2 > x_1$ , we have

$$\begin{aligned} |\mathfrak{N}(\psi)(x_1) - \mathfrak{N}(\psi)(x_2)| &\leq \|\mathfrak{C}\left(\frac{x_2^c}{c}\right) - \mathfrak{C}\left(\frac{x_1^c}{c}\right)\|_{B(\mathcal{F})} \|\eta\| \\ &\quad + \|\mathfrak{S}\left(\frac{x_2^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c}{c}\right)\|_{B(\mathcal{F})} \|\vartheta\| \\ &\quad + \int_0^{x_1} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c-x^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c-x^c}{c}\right)\|_{B(\mathcal{F})} |\Psi(x, \psi_x)| dx \\ &\quad + \int_{x_1}^{x_2} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c-x^c}{c}\right)\|_{B(\mathcal{F})} |\Psi(x, \psi_x)| dx \\ &\leq \|\mathfrak{C}\left(\frac{x_2^c}{c}\right) - \mathfrak{C}\left(\frac{x_1^c}{c}\right)\|_{B(\mathcal{F})} \|\eta\| \end{aligned}$$

$$\begin{aligned}
 &+ \left\| \mathfrak{S}\left(\frac{x_2^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c}{c}\right) \right\|_{B(\mathcal{F})} \|\vartheta\| \\
 &+ \int_0^{x_1} x^{c-1} \left\| \mathfrak{S}\left(\frac{x_2^c - x^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c - x^c}{c}\right) \right\|_{B(\mathcal{F})} |\Psi(x, \psi_x)| dx \\
 &+ \mathcal{M}_c^{\mathfrak{S}} \int_{x_1}^{x_2} x^{c-1} |\Psi(x, \psi_x)| dx.
 \end{aligned}$$

As  $x_1 \rightarrow x_2$ , the uniform continuity property of the operators  $\mathfrak{C}(s)$  and  $\mathfrak{S}(s)$  indicate that the right part of the previous inequality converges to zero. This confirms the equicontinuity of  $\mathfrak{N}$ .

Additionally, we establish the equiconvergence of  $\mathfrak{N}(B_I)$ . For  $s \in I$  and  $\psi \in B_I$ , we find

$$|\mathfrak{N}(\psi)(s)| \leq \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} l \int_0^s x^{c-1} \mathcal{O}(x) dx.$$

Consequently,

$$|\mathfrak{N}(\psi)(s)| \rightarrow l', \text{ as } s \rightarrow +\infty,$$

where  $l' \leq \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} l \mathcal{O}^*$ . Here  $\mathcal{O}^* := \sup_{s \in I} \int_0^s x^{c-1} \mathcal{O}(x) dx$ . Therefore,

$$|\mathfrak{N}(\psi)(s) - \mathfrak{N}(\psi)(+\infty)| \rightarrow 0, s \rightarrow +\infty.$$

Lastly, we prove that the condition of Meir-Keeler’s theorem is satisfied. For any given  $\beta > 0$ , we demonstrate the existence of a  $\mu > 0$  such that:

$$\beta < \mathfrak{H}_I(\Delta) < \beta + \mu \Rightarrow \mathfrak{H}_I(\mathfrak{N}(\Delta)) \leq \beta, \text{ for any } \Delta \subset B_I,$$

where  $\mathfrak{H}_I(\Delta) = \sup_{s \in I} \mathfrak{H}(\Delta(s))$ .

By using the characteristics of the measure of noncompactness (see eg. [12, 21, 35]) and the associated Lemma 2.5 [33], we get

$$\begin{aligned}
 \mathfrak{H}(\mathfrak{N}(\Delta)(s)) &\leq \mathcal{M}_c^{\mathfrak{S}} \int_0^s x^{c-1} \mathcal{O}(x) \mathfrak{H}(\Delta(x)) dx \\
 &\leq \mathcal{M}_c^{\mathfrak{S}} \int_0^s x^{c-1} \mathcal{O}(x) dx \mathfrak{H}_I(\Delta) \\
 &\leq \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* \mathfrak{H}_I(\Delta).
 \end{aligned}$$

Since  $\mathfrak{N}(\Delta)$  is bounded and equicontinuous for all  $\Delta \subset B_I$  then:

$$\mathfrak{H}_I(\mathfrak{N}(\Delta)) = \sup_{s \in I} \mathfrak{H}(\mathfrak{N}(\Delta)(s)).$$

Therefore,  $\mathfrak{H}_I(\mathfrak{N}(\Delta)) \leq \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* \mathfrak{H}_I(\Delta) \leq \beta \Rightarrow \mathfrak{H}_I(\Delta) \leq \frac{\beta}{\mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^*}$ . Then, for given  $\beta > 0$  and taking  $\mu = \left(\frac{1 - \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^*}{\mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^*}\right) \beta - \epsilon$  such that  $\epsilon > 0$ , we obtain:

$$\beta < \mathfrak{H}_I(\Delta) < \beta + \mu \Rightarrow \mathfrak{H}_I(\mathfrak{N}(\Delta)) \leq \beta, \text{ for any } \Delta \subset B_I.$$

We conclude that  $\mathfrak{N}$  is a Meir-Keeler operator with condensing properties.

Through these steps, we ensure that the conditions required for Meir-Keeler’s fixed-point Theorem 2.6 are satisfied by  $\mathfrak{N} : B_I \rightarrow B_I$ . Therefore, we may conclude that  $\mathfrak{N}$  has a fixed point  $\psi$  that provides a mild solution to the problem (1) – (2).

3.2. Uniqueness results

We now introduce our primary result on the existence and uniqueness of solutions for problem (1) – (2). Before demonstrating this result, we first present the following condition.

(iii) The given function  $\Psi : I \times C(\mathcal{H}, \mathcal{F}) \rightarrow \mathcal{F}$  is carathéodory function and there exist a continuous function  $\mathcal{O} : I \rightarrow I$  that satisfies

$$|\Psi(s, u) - \Psi(s, v)| \leq \mathcal{O}(s)\|u - v\|,$$

$$\Psi^* = \sup_{s \in I} \int_0^s x^{c-1} \Psi(x, 0) dx < \infty,$$

$$\mathfrak{H}(\Psi(s, \mathcal{D})) \leq \mathcal{O}(s)\mathfrak{H}(\mathcal{D}),$$

and  $\mathcal{O}^* := \sup_{s \in I} \int_0^s x^{c-1} \mathcal{O}(x) dx < \infty$ , for all  $s \in I$ ,  $u, v \in C(\mathcal{H}, \mathcal{F})$ , bounded set  $\mathcal{D} \subset C([-a, +\infty), \mathcal{F})$  and  $0 < c \leq 1$ .

**Theorem 3.3.** Under the conditions (ii) and (iii) being satisfied, if the product  $\mathcal{M}_c^\xi \mathcal{O}^* < 1$ , then the system defined by equations (1) – (2) possesses a unique mild solution over the space  $BC([-a, +\infty), \mathcal{F})$ .

**Proof:** We observe that following analogous steps to those in the proof of Theorem 3.2, we establish the existence of a mild solution with the modified radius estimate  $l \geq \frac{\mathcal{M}_c^\xi \|\eta\| + \mathcal{M}_c^\xi \|\vartheta\| + \mathcal{M}_c^\xi \Psi^*}{1 - \mathcal{M}_c^\xi \mathcal{O}^*}$ .

Next, we establish uniqueness: Assume  $\psi_1$  and  $\psi_2$  are two mild solutions to problem (1) – (2). Then,

$$\begin{aligned} |\psi_1(s) - \psi_2(s)| &= |\Re\psi_1(s) - \Re\psi_2(s)| \\ &\leq \mathcal{M}_c^\xi \int_0^s x^{c-1} |\Psi(x, \psi_{1x}) - \Psi(x, \psi_{2x})| dx \\ &\leq \mathcal{M}_c^\xi \int_0^s x^{c-1} \mathcal{O}(x) |\psi_{1x} - \psi_{2x}| dx \\ &\leq \mathcal{M}_c^\xi \mathcal{O}^* \|\psi_1 - \psi_2\|_{BC}. \end{aligned}$$

This implies  $(1 - \mathcal{M}_c^\xi \mathcal{O}^*) \|\psi_1 - \psi_2\|_{BC} \leq 0$ , and consequently  $\psi_1 = \psi_2$ . Thus, we have established the uniqueness of the mild solution.

3.3. Attractiveness of Mild Solutions

In this section, we explore the local attractiveness of solutions to problem (1)-(2).

**Theorem 3.4.** Assuming conditions (ii) and (iii) are satisfied, and if  $\mathcal{M}_c^\xi \mathcal{O}^* < 1$ , let  $\psi^*$  be a solution of (1) – (2), and  $\bar{B}(\psi^*, \tau)$  be the closed ball in  $BC$  such that

$$\tau \geq \frac{\mathcal{M}_c^\xi \|\eta\| + \mathcal{M}_c^\xi \|\vartheta\| + \mathcal{M}_c^\xi \Psi^*}{1 - \mathcal{M}_c^\xi \mathcal{O}^*}.$$

Then, the problem (1)-(2) exhibits attractiveness.

**Proof.** Considering  $\psi \in \bar{B}(\psi^*, \tau)$ , utilizing (ii) and (iii), we obtain:

$$\begin{aligned} |\Re(\psi)(s) - \psi^*(s)| &= |\Re(\psi)(s) - \Re(\psi^*)(s)| \\ &\leq \mathcal{M}_c^\xi \int_0^s x^{c-1} |\Psi(x, \psi_x^*) - \Psi(x, \psi_x)| dx \\ &\leq \mathcal{M}_c^\xi \int_0^s x^{c-1} \mathcal{O}(x) \|\psi_x^* - \psi_x\| dx \end{aligned}$$

$$\begin{aligned} &\leq \mathcal{M}_c^\ominus \mathcal{O}^* \|\psi^* - \psi\|_{BC} \\ &\leq \mathcal{M}_c^\ominus \mathcal{O}^* \tau \\ &\leq \tau. \end{aligned}$$

Consequently,  $\mathfrak{N}(\overline{B}(\psi^*, \tau)) \subset \overline{B}(\psi^*, \tau)$ . Thus, for any solutions  $\psi, \tilde{\psi} \in \overline{B}(\psi^*, \tau)$  of (1) – (2) and  $s \in I$ , we have

$$|\psi(s) - \tilde{\psi}(s)| \leq \mathcal{M}_c^\ominus \mathcal{O}^* \|\tilde{\psi} - \psi\|_{BC},$$

then

$$(1 - \mathcal{M}_c^\ominus \mathcal{O}^*) \|\tilde{\psi} - \psi\|_{BC} \leq 0.$$

Hence

$$\|\tilde{\psi} - \psi\|_{BC} = 0.$$

As a result, the problem solutions (1) – (2) are locally attractive.

### 3.4. Controllability results

In this section, we present a controllability result for the system (3)-(4). Prior to that, we introduce a specific type of solutions for problem (3)-(4).

**Definition 3.5.** We define the mild solution  $\psi \in C([-a, +\infty), \mathcal{F})$  of the problem (3) – (4) as follows

$$\psi(s) = \begin{cases} \eta(s), & \text{if } s \in \mathcal{H}; \\ \mathfrak{C}\left(\frac{s^c}{c}\right) \eta(0) + \mathfrak{S}\left(\frac{s^c}{c}\right) \vartheta + \int_0^s x^{c-1} \mathfrak{S}\left(\frac{s^c-x^c}{c}\right) \mathcal{B}\mathcal{U}(x)dx \\ \quad + \int_0^s x^{c-1} \mathfrak{S}\left(\frac{s^c-x^c}{c}\right) \Psi(x, \psi_x) dx, & \text{if } s \in I. \end{cases}$$

**Definition 3.6.** The system (3) – (4) is considered controllable if, for every initial function  $\eta \in C(\mathcal{H}, \mathcal{F})$  and  $\hat{\psi} \in \mathcal{F}$ , there exists some  $n \in \mathbb{N}$  and some control  $\mathcal{U} \in L^2([0, n], \mathcal{F})$  such that the mild solution  $\psi(\cdot)$  of this problem satisfies the terminal condition  $\psi(n) = \hat{\psi}$ .

We will consider the assumptions (i) – (ii) from Section 3, and we introduce the following additional assumptions:

(iv) For all  $n$  integer, the linear operator  $\mathfrak{B} : L^2([0, n], \mathcal{F}) \rightarrow \mathcal{F}$  defined by

$$\mathfrak{B}\mathcal{U} = \int_0^n x^{c-1} \mathfrak{S}\left(\frac{n^c-x^c}{c}\right) \mathcal{B}\mathcal{U}(x)dx,$$

possesses a pseudo-invertible operator  $\tilde{\mathfrak{B}}^{-1}$ , which maps functions from  $L^2([0, n], \mathcal{F})$  to the space  $L^2([0, n], \mathcal{F})$  excluding the kernel of  $\mathfrak{B}$ , and is bounded. Additionally,  $\mathcal{B}$  is bounded, satisfying:

$$\|\mathcal{B}\| \leq \tilde{\mathcal{N}} \text{ and } \|\tilde{\mathfrak{B}}^{-1}\| \leq \tilde{\mathcal{N}}_1.$$

(v) There exists a continuous function  $\mathcal{K}_{\mathfrak{B}} : [0, n] \rightarrow \mathbb{R}_+$  such that: for any bounded subset  $\mathcal{D} \subset \mathcal{F}$ , we have :  $\mathfrak{S}(\tilde{\mathfrak{B}}^{-1}(\mathcal{D}))(s) \leq \mathcal{K}_{\mathfrak{B}}(s)\mathfrak{S}(\mathcal{D}), s \in I$  and  $\mathcal{K}' := \sup_{s \in I} \int_0^s x^{c-1} \mathcal{K}_{\mathfrak{B}}(x)dx < \infty$  for all  $0 < c \leq 1$ .

**Theorem 3.7.** Assuming that (i) – (ii) and (iv) – (v) hold. If

$$\max\{\mathcal{M}_c^\ominus [\mathcal{O}^* + \tilde{\mathcal{N}}\tilde{\mathcal{N}}_1 \frac{n^c}{c} (1 + \mathcal{M}_c^\ominus \mathcal{O}^*)], \mathcal{M}_c^\ominus \mathcal{O}^* (1 + \mathcal{M}_c^\ominus \tilde{\mathcal{N}} \mathcal{K}')\} < 1,$$

then the problem (3) – (4) is controllable on  $[-a, +\infty)$ .



**Proof.** We reformulate the problem (3) – (4) into a fixed-point issue. Define the operator  $\mathfrak{R} : BC_\infty \rightarrow BC_\infty$  as follows:

$$\mathfrak{R}(\psi)(s) = \begin{cases} \eta(s), & \text{if } s \in \mathcal{H}; \\ \mathfrak{C}\left(\frac{s^c}{c}\right) \eta(0) + \mathfrak{S}\left(\frac{s^c}{c}\right) \vartheta + \int_0^s x^{c-1} \mathfrak{S}\left(\frac{s^c-x^c}{c}\right) \mathcal{B}\mathcal{U}(x) dx \\ + \int_0^s x^{c-1} \mathfrak{S}\left(\frac{s^c-x^c}{c}\right) \Psi(x, \psi_x) dx, & \text{if } s \in [0, +\infty). \end{cases}$$

Utilizing assumption (iv), we can define the control for any arbitrary function  $\psi(\cdot)$ :

$$\mathcal{U}_\psi(s) = \tilde{\mathfrak{B}}^{-1} \left[ \hat{\psi} - \mathfrak{C}\left(\frac{n^c}{c}\right) \eta(0) - \mathfrak{S}\left(\frac{n^c}{c}\right) \vartheta - \int_0^n x^{c-1} \mathfrak{S}\left(\frac{n^c-x^c}{c}\right) \Psi(x, \psi_x) dx \right](s).$$

Such that, we have

$$\begin{aligned} |\mathcal{U}_\psi(s)| &\leq \|\tilde{\mathfrak{B}}^{-1}\| \left[ |\hat{\psi}| + \mathcal{M}_c^{\mathfrak{C}} |\eta(0)| + \mathcal{M}_c^{\mathfrak{S}} |\vartheta| + \mathcal{M}_c^{\mathfrak{S}} \int_0^n x^{c-1} \mathcal{O}(x) |\psi_x| dx \right] \\ &\leq \tilde{\mathcal{N}}_1 \left[ |\hat{\psi}| + \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* \|\psi\|_{BC_\infty} \right] \end{aligned}$$

The operator  $\mathfrak{R}$  maps  $BC_\infty$  into  $BC_\infty$ . Specifically, the mapping  $\mathfrak{R}(\psi)$  is continuous on  $[-a, n]$  for any  $\psi \in BC_\infty$  we have:

$$\begin{aligned} |\mathfrak{R}(\psi)(s)| &\leq \|\mathfrak{C}\left(\frac{s^c}{c}\right)\| |\eta(0)| + \|\mathfrak{S}\left(\frac{s^c}{c}\right)\| |\vartheta(0)| + \int_0^s x^{c-1} \|\mathfrak{S}\left(\frac{s^c-x^c}{c}\right)\| |\Psi(x, \psi_x)| dx \\ &\quad + \int_0^s x^{c-1} \|\mathfrak{S}\left(\frac{s^c-x^c}{c}\right)\| \|\mathcal{B}\| |\mathcal{U}_\psi(x)| dx \\ &\leq \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \int_0^s x^{c-1} \mathcal{O}(x) \|\psi_x\| dx \\ &\quad + \mathcal{M}_c^{\mathfrak{S}} \int_0^s x^{c-1} \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \left[ |\hat{\psi}| + \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* \|\psi\|_{BC_\infty} \right] dx \\ &\leq \mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* \|\psi\|_{BC_\infty} + \mathcal{M}_c^{\mathfrak{S}} \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \frac{n^c}{c} \left[ |\hat{\psi}| + \mathcal{M}_c^{\mathfrak{C}} \|\eta\| \right. \\ &\quad \left. + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\| + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^* \|\psi\|_{BC_\infty} \right] \\ &\leq (\mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\|) (1 + \mathcal{M}_c^{\mathfrak{S}} \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \frac{n^c}{c}) + \mathcal{M}_c^{\mathfrak{S}} \|\psi\|_{BC_\infty} \left[ \mathcal{O}^* + \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \frac{n^c}{c} (1 + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^*) \right]. \end{aligned}$$

So,  $\mathfrak{R} \in BC_\infty$ .

Furthermore, suppose  $l \geq \frac{(\mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\|) (1 + \mathcal{M}_c^{\mathfrak{S}} \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \frac{n^c}{c})}{1 - \mathcal{M}_c^{\mathfrak{S}} \left[ \mathcal{O}^* + \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \frac{n^c}{c} (1 + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^*) \right]}$ , and let  $B_l$  denote the closed ball in  $BC_\infty$  centered at the origin with radius  $l$ . Let  $\psi \in B_l$  and  $s \in I$ , we get

$$|\mathfrak{R}(\psi)(s)| \leq (\mathcal{M}_c^{\mathfrak{C}} \|\eta\| + \mathcal{M}_c^{\mathfrak{S}} \|\vartheta\|) (1 + \mathcal{M}_c^{\mathfrak{S}} \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \frac{n^c}{c}) + \mathcal{M}_c^{\mathfrak{S}} l \left[ \mathcal{O}^* + \tilde{\mathcal{N}} \tilde{\mathcal{N}}_1 \frac{n^c}{c} (1 + \mathcal{M}_c^{\mathfrak{S}} \mathcal{O}^*) \right].$$

Thus, we get that

$$\|\mathfrak{R}(\psi)\|_{BC_\infty} \leq l.$$

We embark on the validation process to ensure that  $\mathfrak{R} : B_l \rightarrow B_l$  satisfies the criteria outlined in Meir Keeler’s fixed-point Theorem 2.6.

Firstly, we ascertain the continuity of  $\mathfrak{N}$  within  $B_I$ . Let  $\{\psi_k\}$  be a sequence such that  $\psi_k \rightarrow \psi$  in  $B_I$ . We observe that:

$$\begin{aligned} |\mathfrak{N}(\psi_k)(s) - \mathfrak{N}(\psi)(s)| &\leq \mathcal{M}_c^\ominus \int_0^s x^{c-1} |\Psi(x, (\psi_x)_k) - \Psi(x, \psi_x)| dx \\ &\quad + \mathcal{M}_c^\ominus \widetilde{\mathcal{N}} \int_0^s x^{c-1} |\mathcal{U}_{\psi_k}(x) - \mathcal{U}_\psi(x)| dx \\ &\leq \mathcal{M}_c^\ominus \int_0^s x^{c-1} |\Psi(x, (\psi_x)_k) - \Psi(x, \psi_x)| dx \\ &\quad + \mathcal{M}_c^\ominus \widetilde{\mathcal{N}} \widetilde{\mathcal{N}}_1 \int_0^s x^{c-1} \left[ |\widehat{\psi}_k - \widehat{\psi}| \right. \\ &\quad \left. + \mathcal{M}_c^\ominus \int_0^n \tau^{c-1} |\Psi(\tau, (\psi_\tau)_k) - \Psi(\tau, \psi_\tau)| d\tau \right] dx \\ &\leq \mathcal{M}_c^\ominus \left( 1 + \mathcal{M}_c^\ominus \widetilde{\mathcal{N}} \widetilde{\mathcal{N}}_1 \frac{n^c}{c} \right) \int_0^n x^{c-1} |\Psi(x, (\psi_x)_k) - \Psi(x, \psi_x)| dx \\ &\quad + \mathcal{M}_c^\ominus \widetilde{\mathcal{N}} \widetilde{\mathcal{N}}_1 \frac{n^c}{c} |\widehat{\psi}_k - \widehat{\psi}|. \end{aligned}$$

According to (i), we have  $\Psi(x, (\psi_x)_k) \rightarrow \Psi(x, \psi_x)$  as  $k \rightarrow +\infty$  for almost every  $x \in [0, n]$ . Application of the Lebesgue dominated convergence Theorem yields:

$$\|\mathfrak{N}(\psi_k) - \mathfrak{N}(\psi)\|_{BC_\infty} \rightarrow 0, \text{ as } k \rightarrow \infty.$$

Consequently,  $\mathfrak{N}$  exhibits continuity.

So by the previous steps we can conclude that  $\mathfrak{N}(B_I) \subset B_I$ .

Next, we confirm the equicontinuity of  $\mathfrak{N}(B_I)$  over every compact interval  $J = [0, n]$ . let  $x_1, x_2 \in J$  with  $x_2 > x_1$  we have

$$\begin{aligned} |\mathfrak{N}(\psi)(x_1) - \mathfrak{N}(\psi)(x_2)| &\leq \|\mathfrak{C}\left(\frac{x_2^c}{c}\right) - \mathfrak{C}\left(\frac{x_1^c}{c}\right)\|_{B(\mathcal{F})} \|\eta\| + \|\mathfrak{S}\left(\frac{x_2^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c}{c}\right)\|_{B(\mathcal{F})} \|\vartheta\| \\ &\quad + \int_0^{x_1} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c - x^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c - x^c}{c}\right)\|_{B(\mathcal{F})} |\Psi(x, \psi_x)| dx \\ &\quad + \int_{x_1}^{x_2} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c - x^c}{c}\right)\|_{B(\mathcal{F})} |\Psi(x, \psi_x)| dx \\ &\quad + \int_0^{x_1} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c - x^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c - x^c}{c}\right)\|_{B(\mathcal{F})} \|\mathcal{B}\| \|\mathcal{U}_\psi(x)\| dx \\ &\quad + \int_{x_1}^{x_2} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c - x^c}{c}\right)\|_{B(\mathcal{F})} \|\mathcal{B}\| \|\mathcal{U}_\psi(x)\| dx \\ &\leq \|\mathfrak{C}\left(\frac{x_2^c}{c}\right) - \mathfrak{C}\left(\frac{x_1^c}{c}\right)\|_{B(\mathcal{F})} \|\eta\| + \|\mathfrak{S}\left(\frac{x_2^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c}{c}\right)\|_{B(\mathcal{F})} \|\vartheta\| \\ &\quad + \int_0^{x_1} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c - x^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c - x^c}{c}\right)\|_{B(\mathcal{F})} |\Psi(x, \psi_x)| dx \\ &\quad + \mathcal{M}_c^\ominus \int_{x_1}^{x_2} x^{c-1} |\Psi(x, \psi_x)| dx \\ &\quad + \int_0^{x_1} x^{c-1} \|\mathfrak{S}\left(\frac{x_2^c - x^c}{c}\right) - \mathfrak{S}\left(\frac{x_1^c - x^c}{c}\right)\|_{B(\mathcal{F})} \|\mathcal{B}\| \|\mathcal{U}_\psi(x)\| dx \\ &\quad + \mathcal{M}_c^\ominus \widetilde{\mathcal{N}} \int_{x_1}^{x_2} x^{c-1} \|\mathcal{U}_\psi(x)\| dx. \end{aligned}$$

As  $x_1 \rightarrow x_2$ , the uniformly continuity property of  $\mathfrak{C}(s)$  and  $\mathfrak{S}(s)$  indicate that the right part of the

previous inequality converges to zero. This confirms the equicontinuity of  $\mathfrak{R}$ .

Additionally, we establish the equiconvergence of  $\mathfrak{R}(B_I)$ . For  $s \in J$  and  $\psi \in B_I$ , we find

$$|\mathfrak{R}(\psi)(s)| \leq (\mathcal{M}_c^{\mathfrak{C}}\|\eta\| + \mathcal{M}_c^{\mathfrak{E}}\|\vartheta\|)(1 + \mathcal{M}_c^{\mathfrak{E}}\widetilde{\mathcal{N}}\widetilde{\mathcal{N}}_1\frac{n^c}{c}) + \mathcal{M}_c^{\mathfrak{E}}I \left[ \int_0^s x^{c-1}\mathcal{O}(x)dx + \widetilde{\mathcal{N}}\widetilde{\mathcal{N}}_1\frac{n^c}{c}(1 + \mathcal{M}_c^{\mathfrak{E}}\int_0^s x^{c-1}\mathcal{O}(x)dx) \right].$$

Consequently,

$$|\mathfrak{R}(\psi)(s)| \rightarrow l', \text{ as } s \rightarrow +\infty.$$

where  $l' \leq (\mathcal{M}_c^{\mathfrak{C}}\|\eta\| + \mathcal{M}_c^{\mathfrak{E}}\|\vartheta\|)(1 + \mathcal{M}_c^{\mathfrak{E}}\widetilde{\mathcal{N}}\widetilde{\mathcal{N}}_1\frac{n^c}{c}) + \mathcal{M}_c^{\mathfrak{E}}I \left[ \mathcal{O}^* + \widetilde{\mathcal{N}}\widetilde{\mathcal{N}}_1\frac{n^c}{c}(1 + \mathcal{M}_c^{\mathfrak{E}}\mathcal{O}^*) \right]$ . Here  $\mathcal{O}^* := \sup_{s \in I} \int_0^s x^{c-1}\mathcal{O}(x)dx$ . Therefore,

$$|\mathfrak{R}(\psi)(s) - \mathfrak{R}(\psi)(+\infty)| \rightarrow 0, s \rightarrow +\infty.$$

Lastly, we validate the satisfaction of the Meir-Keeler’s type condition. For all  $\beta > 0$  given. we prove that there exists  $\mu > 0$  such that:

$$\beta < \mathfrak{H}_I(\Delta) < \beta + \mu \Rightarrow \mathfrak{H}_I(\mathfrak{R}(\Delta)) \leq \beta, \text{ for any } \Delta \subset B_I,$$

where

$$\mathfrak{H}_I(\Delta) = \sup_{s \in I} \mathfrak{H}(\Delta(s)).$$

By using the properties of the measure of noncompactness (see eg. [12, 21, 35]) and Lemma 2.5 [33], we get

$$\mathfrak{H}(\mathfrak{R}(\Delta)(s)) \leq \mathcal{M}_c^{\mathfrak{E}}\int_0^s x^{c-1}\mathcal{O}(x)\mathfrak{H}(\Delta(x))dx + \mathcal{M}_c^{\mathfrak{E}}\widetilde{\mathcal{N}}\int_0^s x^{c-1}\mathfrak{H}(\mathcal{U}_\Delta(x))dx.$$

We have,

$$\begin{aligned} \mathfrak{H}(\mathcal{U}_\Delta(x)) &\leq \mathcal{K}_{\mathfrak{B}}(x)\mathcal{M}_c^{\mathfrak{E}}\int_0^n \tau^{c-1}\mathfrak{H}(\Psi(\tau, \Delta(\tau)))d\tau \\ &\leq \mathcal{K}_{\mathfrak{B}}(x)\mathcal{M}_c^{\mathfrak{E}}\int_0^n \tau^{c-1}\mathcal{O}(\tau)\mathfrak{H}(\Delta(\tau))d\tau \\ &\leq \mathcal{K}_{\mathfrak{B}}(x)\mathcal{M}_c^{\mathfrak{E}}\mathcal{O}^*\mathfrak{H}_I(\Delta), \end{aligned}$$

which implies

$$\begin{aligned} \mathfrak{H}(\mathfrak{R}(\Delta)(s)) &\leq \mathcal{M}_c^{\mathfrak{E}}\int_0^s x^{c-1}\mathcal{O}(x)dx\mathfrak{H}_I(\Delta) + \mathcal{M}_c^{\mathfrak{E}}\widetilde{\mathcal{N}}\int_0^s x^{c-1}\mathcal{K}_{\mathfrak{B}}(x)\mathcal{M}_c^{\mathfrak{E}}\mathcal{O}^*\mathfrak{H}_I(\Delta)ds \\ &\leq \mathcal{M}_c^{\mathfrak{E}}\mathcal{O}^*\mathfrak{H}_I(\Delta) + \mathcal{M}_c^{\mathfrak{E}^2}\widetilde{\mathcal{N}}\mathcal{O}^*\mathcal{K}'\mathfrak{H}_I(\Delta) \\ &\leq (\mathcal{M}_c^{\mathfrak{E}}\mathcal{O}^* + \mathcal{M}_c^{\mathfrak{E}^2}\widetilde{\mathcal{N}}\mathcal{O}^*\mathcal{K}')\mathfrak{H}_I(\Delta). \end{aligned}$$

Since  $\mathfrak{R}(\Delta)$  is bounded and equicontinuous for all  $\Delta \subset B_I$  then:

$$\mathfrak{H}_I(\mathfrak{R}(\Delta)) = \sup_{s \in I} \mathfrak{H}(\mathfrak{R}(\Delta)(s)).$$

Therefore,  $\mathfrak{H}_I(\mathfrak{R}(\Delta)) \leq (\mathcal{M}_c^{\mathfrak{E}}\mathcal{O}^* + \mathcal{M}_c^{\mathfrak{E}^2}\widetilde{\mathcal{N}}\mathcal{O}^*\mathcal{K}')\mathfrak{H}_I(\Delta) \leq \beta \Rightarrow \mathfrak{H}_I(\Delta) \leq \frac{\beta}{\mathcal{M}_c^{\mathfrak{E}}\mathcal{O}^* + \mathcal{M}_c^{\mathfrak{E}^2}\widetilde{\mathcal{N}}\mathcal{O}^*\mathcal{K}'}$ . Then, for

given  $\beta > 0$  and taking  $\mu = \left( \frac{1 - (\mathcal{M}_\epsilon^{\tilde{\mathcal{O}}^*} + \mathcal{M}_\epsilon^{\tilde{\mathcal{N}}^2} \tilde{\mathcal{N}} \mathcal{O}^* \mathcal{K}')}{\mathcal{M}_\epsilon^{\tilde{\mathcal{O}}^*} + \mathcal{M}_\epsilon^{\tilde{\mathcal{N}}^2} \tilde{\mathcal{N}} \mathcal{O}^* \mathcal{K}'} \right) \beta - \epsilon$  such that  $\epsilon > 0$ , we obtain:

$$\beta < \mathfrak{H}_I(\Delta) < \beta + \mu \Rightarrow \mathfrak{H}_I(\mathfrak{N}(\Delta)) \leq \beta, \text{ for any } \Delta \subset B_I.$$

Hence  $\mathfrak{N}$  is a Meir-Keeler condensing operator.

Through these steps, we ensure that the conditions required for Meir-Keeler’s fixed-point Theorem 2.6 are satisfied by  $\mathfrak{N} : B_I \rightarrow B_I$ . Therefore, we may conclude that  $\mathfrak{N}$  has a fixed point  $\psi$  that provides a mild solution to the problem (3) – (4).

#### 4. Examples

**Example 4.1.** To showcase the practical application of our results, let  $\mathcal{E}$  denote a nonempty bounded open set in  $\mathbb{R}^2$ . We explore the following conformable fractional differential equation:

$$D_s^{\frac{1}{3}} [D_s^{\frac{1}{3}} \psi(s, x)] = D_x^2 \psi(s, x) + \Psi(s, \psi(s - a, x)), \quad x \in \mathcal{E}, \quad s \in [0, +\infty); \tag{5}$$

$$\psi(s, x) = 0, \quad s \in [0, +\infty), \quad x \in \partial \mathcal{E}; \tag{6}$$

$$\psi(s, x) = \eta(s, x); \quad D_s^{\frac{1}{3}} [\psi(0, x)] = \vartheta, \quad s \in [-a, 0], \quad x \in \mathcal{E}. \tag{7}$$

Here,  $a > 0$  and we have

$$\Psi(s, \psi(s - a, x)) = \frac{\exp -s}{7} \sin \psi(s - a, x).$$

Taking  $\mathcal{F} = L^2(\mathcal{E})$  and defining  $\mathfrak{B}$  as follows:

$$D(\mathfrak{B}) = \{ \psi \in \mathcal{H}(\mathcal{F}), \psi(x)|_{x \in \partial \mathcal{E}} = 0 \},$$

$$\mathfrak{B}\psi = D_x^2 \psi, \quad \psi \in D(\mathfrak{B}).$$

It is well known the operator  $\mathfrak{B}$  generates a cosine family  $((\mathfrak{C}(s))_{s \in \mathbb{R}}, (\mathfrak{S}(s))_{s \in \mathbb{R}})$ . Additionally, it follows that

$$\|\mathfrak{C}(s)\| \leq 1 \text{ and } \|\mathfrak{S}(s)\| \leq 1, \text{ for all } s \in [0, +\infty).$$

Thus, to apply our theorems on existence and attractivity, we require  $\mathcal{O}^* < 1$ .

The function  $\Psi(s, \psi(s - a, x)) = \frac{\exp -s}{7} \sin \psi(s - a, x)$  is Carathéodory and

$$|\Psi(s, \psi_1(s - a, x)) - \Psi(s, \psi_2(s - a, x))| \leq \frac{\exp -s}{7} |\psi_1(s - a, x) - \psi_2(s - a, x)|,$$

thus  $\mathcal{O}(s) = \frac{\exp -s}{7}$ . Moreover, we have

$$\mathcal{O}^* = \sup \left\{ \int_0^s x^{-\frac{2}{3}} \frac{\exp -x}{7} dx, \quad s \in [0, +\infty) \right\} = \frac{\Gamma(\frac{1}{3})}{7} \simeq 0.3827 < 1, \quad \Psi_0 = 0.$$

Then, by [21, 31], the problem (1)-(3) is an abstract formulation of the problem (5)-(7), and conditions (i) – (iii) are satisfied. Theorem 3.3 implies that the problem (5)-(7) has a unique mild solution on BC, which is attractive by Theorem 3.4.

**Example 4.2.** To showcase the practical application of our results, let  $\mathcal{E}$  denote a nonempty bounded open set in  $\mathbb{R}^2$ . We explore the following conformable fractional differential equation:

$$D_s^{\frac{1}{3}} [D_s^{\frac{1}{3}} \psi(s, x)] = D_x^2 \psi(s, x) + \Psi(s, \psi(s - a, x)) + g(x)\mathcal{U}(s), \quad x \in \mathcal{E}, \quad s \in [0, +\infty); \tag{8}$$

$$\psi(s, x) = 0, s \in [0, +\infty), x \in \partial\mathcal{E}; \quad (9)$$

$$\psi(s, x) = \eta(s, x); D_s^{\frac{1}{3}}[\psi(0, x)] = \vartheta, s \in [-a, 0], x \in \mathcal{E}. \quad (10)$$

Here,  $\mathfrak{P}$ ,  $\Psi$ ,  $\mathcal{Y}$  and  $\eta$  are as in problem (5) – (6),  $g : \mathcal{E} \rightarrow \mathcal{F}$  is a continuous function and  $\mathcal{U} : I \rightarrow \mathcal{F}$  is a given control.

Then, by [21, 31], the problem (3)-(4) is an abstract formulation of the problem (8)-(10), and conditions (i) – (iii) are satisfied. Theorem 3.7 implies that the problem (8)-(10) is controllable on  $[-a, +\infty)$ .

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